## WAVEGUIDE MODEL OF WIRELESS COMMUNICATION CHANNELS INSIDE BUILDINGS

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#### Abstract

A plenty of wireless data transmission networks now works inside of buildings.

Specifics of electromagnetic waves propagation of in these conditions are:

Presence of possible obstacles in a propagation path of a signal;

Multiple propagation paths of a signal because of numerous reflections.

The models communication channels offered at present time inside of buildings do not fully consider these specifics therefore do not possess satisfactory accuracy of calculation. Besides disadvantage of these models is their "instability" to amount of initial data. Initial data for creation of mathematical model of radio channels inside buildings is the plan of this building. The "steady" model allows to estimate roughly a level of a signal using the minimal amount of initial data about a building or city and leads to improvement of accuracy in the process of data specification. Models applied at the present moment start to work only after initial data about a lay-out of a building or city are full enough and these models do not give substantial improvement of accuracy at their updating.

The analysis of existing radio waves propagation models is performed in this paper.

The task of signal transmission between transmitter and receiver located in the random points inside the building is resolved.

New developed mathematical model of radio waves propagation inside buildings is called wave model.

### Keywords: Mathematical model, propagation, radio channels.

#### **Presenting Author's biography**

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#### 1 Methods of modeling of communication channels inside of buildings

The models of signals propagation inside of buildings offered at the present time can be divided into 3 groups:

Statistical models; not demanding the detailed information on a building except for the general description of its type: an industrial building, hotel, hospital, shopping center, a building of old construction, etc. [1-2].

Empirical single- or multibeam models; based on the analysis of one or several beams connecting transmitting and receiving antennas, for an estimation of a level of an received signal [1], [3-5].

Beam models; in which kvasioptic representation of signals propagation processes is used and reflections from room walls and diffraction at corners [6-8] are considered.

In models of 1-st group losses signal propagation  $L_{\rm p}$  have dependence from distances between antennas d kind

$$L_{p} = L_{p0}(d_{0}) + n \cdot 10 \lg(d / d_{0}), dB$$
(1)

where the exponent n is defined by type of a building.

Such model is applied in a software package [9] and characterized by speed of calculations at which performance it is required to determine only distance between antennas - all other parameters and constants concern to all building as a whole and are set preliminary.

Models of 2-nd group (Montey Keenan [3], Multi-Wall-Model [1]) are based on addition to (1) losses in all walls on a way between receiving and transmitting antennas. In model DPM (Dominant Path Model) [10] concerning the same group to the basic beam the additional beams which are passing through the next rooms in relation to that lay on the main way are added. Thus exact search of points of reflection of a signal is not made.

Models of 3-rd group [6-7] whenever possible fully consider the information on a lay-out of a building. According to them, all probable ways of a signal from the antenna of the transmitter to the antenna of the receiver are defined. For reduction of calculations time some ways of acceleration of computing process [7] are offered.

There are two versions of realization of the given sort of the models, named by trace of beams (ray tracing) and creation of beams (ray launching). The number of considered iterations (reflections, obstacles) depends on computation power of a computer. The majority of models is limited by a maximum by 6-th iterations, including no more than 2 obstacles.

Difractional losses of a signal along each way pay off with use of the geometrical theory of diffraction [10], and factors of reflection - by means of Fresnel formulas [11]. It is possible also to use empirical ratios, calibrated by means of experimental data.

The basic disadvantage of models of 3-rd group is their sensitivity accuracy to the initial data [7-8] as already was marked. In cases when it is inexact data about parameters of walls [11] are specified or to their site, results of calculation essentially worsen [12].

2. waveguide model of radio channels inside of buildings

Let's examine a problem about signaling between located in any points inside of a building a source and the receiver. To receive settlement ratio it is necessary to solve a regional problem about excitation of electromagnetic waves by the set source inside of a building. The decision of this regional problem we shall construct as follows. We shall break internal volume of a building into some final number of blocks so that each of them represented regular structure even along one axis and, probably, non-uniform along two other axes of three-dimensional system of coordinates, perpendicular the first (Fig. 1). Such blocks which we shall name elementary, depending on an internal layout of a building walls of the building, separate rooms or their parts, and also groups of several rooms. Actually, such approach means use of a method of partial areas (MPA) for calculation of an electromagnetic field inside of examined complex structure [13].



Fig. 1 Splitting of a building into elementary blocks

$$\begin{cases} \stackrel{\rho}{E}_{kv}, \stackrel{\rho}{H}_{kv} \end{cases} [13]:$$

$$\begin{cases} \stackrel{\rho}{E}_{k} = \sum_{H} A_{\acute{e}i} \stackrel{\rho}{E}_{\acute{e}i}, \\ \stackrel{\rho}{H}_{k} = \sum_{H} A_{\acute{e}i} \stackrel{\rho}{H}_{\acute{e}i}, \end{cases}$$

$$(2)$$

Having demanded equality of tangential making electromagnetic fields on borders of blocks

$$\begin{cases} P_{k\delta}^{\rho} = E_{m\delta}^{\rho} \\ H_{k\delta}^{\rho} = H_{m\delta}^{\rho}, \end{cases}$$
(3)

we shall receive system of the equations for unknown factors  $A_{\kappa H}$  in Eq. (2).

Generally elementary blocks into which the building is broken, represent structure, regular along one axis and is non-uniform filled along two others, perpendicular the first, axes. Thus, for successful application of an offered method of calculation of characteristics of signals extending inside of buildings it is necessary to investigate characteristics it is non-uniform the filled dielectric structures which commonly can be presented in the form of, represented on fig. 2. Problems of signaling in such structure are in detail examined in [14-15].



Fig. 2. It is non-uniform filled waveguide communication channel

As well as in case of beam models of 3-rd group, application waveguide inside of buildings with a complex lay-out is impossible for model of signals propagation without use of special procedures of acceleration of calculations. Such variant waveguide models is examined in following section.

# 2 The power formulation waveguide models

With the purpose of reception of calculation algorithm of signals propagation characteristics convenient for computer realization inside of buildings we shall examine a problem about excitation of electromagnetic waves in the elementary block - a parallelepiped  $x_1 \times x_2 \times x_3$ ,

Homogeneously filled by not magnetic environment with known size  $s=y/\mu e = tg\partial$  and

Surrounded by walls of infinite thickness with depth of penetration of a field [16]  $\Delta_m$ ,

m=1...,6 Fig.(3). A source we shall consider dot, located in a point of M inside of a parallelepiped.



Fig. 3. The elementary block inside of a building

Field raised in examined structure we shall present in the form of superposition of fields of its own fluctuations as which we shall consider as fashions of homogeneously filled rectangular resonator with losses. Frequencies of the specified fashions are Eq.(4):

$$f_{\nu} = f_{\nu}^{0} \left( 1 + i \left( 2Q_{\nu} \right)^{-1} \right)$$
(4)

where

$$f_{\nu}^{0} = (c/2\pi) ((n\pi/x_{1})^{2} + (l\pi/x_{2})^{2} + (j\pi/x_{3})^{2})^{1/2}$$

- own frequency of v-th fashion of structure for the lack of losses, with c - speed of light,  $Q_v$  - good quality of v-th fashion, and a three of indexes (n, l, j) defines type and, accordingly, the index of v-th fashion appropriated to it. The amplitude Av from which v-th fashion is raised, depends on capacity, type, polarization, the location of a source and inversely proportional differences of squares of frequency  $f_v$  and frequencies of excitation  $f_0$  [17]:

$$A_{v} = b_{v} / \left( f_{v}^{2} - f_{o}^{2} \right)$$
 (5)

where  $b_v$  - a constant depending on an index v.

To examined structure take place repeated reflections of electromagnetic waves from walls. At each reflection the part of energy leaves in a wall, and the remained part dissipates and collects in volume of a room. The size accumulated for the period of exciting fluctuation of energy is equal structure

$$W = \sum_{q} \left| A_{q}^{2} \right| W_{q} \tag{6}$$

where Wq - the electromagnetic energy reserved for the period of fluctuation by q-th fashion with individual amplitude (norm of q-th fashion).

Summation in Eq.(6) should be spent on all fashions of the block. However, presence of quickly decreasing denominator leads to fast convergence Eq.(6) at increase in a difference of frequencies  $|f_v-f_0|$ . For this reason at calculation of the sum of some Eq.(6) it appears sufficient to keep final (though also big enough) number of the fashions belonging set with own frequencies laying in narrow strip  $\Delta$  of frequencies f in a vicinity of frequency  $f_0, \Delta f \leq f_0$ :

$$\Omega = \left\{ f_H : \left| f_H - f_0 \right| \le \Delta f / 2 \right\}.$$

After the termination of transients in structure there comes power balance for averages for the period of exciting fluctuation of powers: power, given by source  $P_{\text{HJIT}}$ , it appears to the equal sum of powers disseminated in volume  $P_0$  and leaving in surrounding walls  $P_{\text{cm}}$ :

$$P_{_{U3\pi}} = P_0 + P_{_{cm}} \tag{7}$$

Power  $P_{cm}$  in Eq.(7) develops of powers  $P_m$  leaving in each of six surrounding structure walls

$$P_{cm} = \sum_{m=1}^{0} P_m$$
 (8)

At  $P_{\mu}/\mu_{\mu} \ll |A_{\mu}^2|W_{\mu}$  entering in Eq.(7) – Eq.(8) powers are determined by following ratio:

$$P_{_{\mathcal{U}3\mathcal{I}}} = \sum_{_{\mathcal{H}\in\Omega}} P_{_{\mathcal{H}}}, \quad P_{_{0}} = \sum_{_{\mathcal{H}\in\Omega}} P_{_{\mathcal{H}0}}, \quad P_{_{m}} = \sum_{_{\mathcal{V}\in\Omega}} P_{_{\mathcal{V}m}}$$
(9)

where  $P_{\nu}$ ,  $P_{\nu 0}$  and  $P_{\nu m}$  - corresponding powers of losses in case of when one is raised in structure at  $\nu$ -th a fashion with amplitude  $A_{\rm H}$ .

Power  $P_{\nu}$ ,  $P_{\nu 0}$  and  $P_{\nu m}$  are defined by means of full  $Q_{\nu}$ , and partial  $Qv_0$  and  $Q_{\nu m}$  good qualities at- $\nu$ -th fluctuation

$$P_{_{H}} = u_{_{H}} |A_{_{H}}^{2}| W_{_{H}} / Q_{_{_{H},}}$$
(10)  
$$P_{_{H}0} = u_{_{H}} |A_{_{H}}^{2}| W_{_{H}} / Q_{_{_{H}0,}}$$
  
$$P_{_{H}m} = u_{_{H}} |A_{_{H}}^{2}| W_{_{H}} / Q_{_{_{H}m,}}$$
  
$$= 2 p f_{_{H},} \qquad 1 / Q_{_{H}} = \sum_{m=0}^{6} (1 / Q_{_{H}m}).$$

In conformity with Eq.(6) – Eq.(10), power of losses in the structure, numerically equal  $P_{_{H3JP}}$  is distributed on set of raised fashions under the law

$$P_{\mu} = \frac{u_{\mu} |A_{\mu}^{2}| W_{\mu} / Q_{\mu}}{\sum_{q} u_{q} |A_{q}^{2}| W_{q} / Q_{q}} P_{u_{3\eta}}$$
(11)

In case of  $\pi_{_{H3R}} \ll x_1$ ,  $x_2$ ,  $x_3$  at run an index q on set of frequencies  $\Omega$  in Eq.(11) takes place fast oscillation constants  $|A^2_{q}|W_{q}/Q_{q}$ . Summation in a denominator Eq. (11) carries out a role of operation of averaging specified oscillations owing to what it appears possible not breaking equality Eq. (7) to consider in Eq. (11) identical sizes  $|A^2_{q}|W_{q}/Q_{q}$  for all  $q \in \Omega$ . Then,

considering, that  $|\mathfrak{U}_q| \approx \mathfrak{U}_{_{H3\Pi}}$  for all fashions from  $\Omega$ , the ratio Eq. (11) becomes:

$$P_{_{H}} = P_{_{\mathcal{U}\mathcal{I}\mathcal{I}}} / \partial_{_{\mathcal{H}}} \tag{12}$$

where

$$\mathcal{D}_{i} = \left\{ \left[ \left( f_{i}^{0} \right)^{2} - \left( f_{\delta \varphi \varphi}^{0} \right)^{2} \right]^{2} + \left( f_{i}^{0} / \mathcal{Q}_{i} \right)^{4} \right\}_{q \in \mathcal{Q}} \left\{ \left[ \left( f_{i}^{0} \right)^{2} - \left( f_{\delta \varphi \varphi}^{0} \right)^{2} \right]^{2} + \left( f_{i}^{0} / \mathcal{Q}_{i} \right)^{4} \right\}^{-1}$$

By means of Eq.(7) - Eq. (12) making powers of losses in examined structure can be written down in the form of

$$P_m = g_m P_{u_{3,n}} \tag{13}$$

where factors  $g_m$  (m=0...,6) are defined by the form of the examined block, character of its filling, characteristics of walls surrounding it, but do not depend on the location and characteristics of an orientation of a source of an electromagnetic field and are equal:

$$g_m = \sum_{i \in \Omega} \frac{Q_i}{\partial_i Q_{i\hat{o}}}$$
(14)

Independence in Eq.(13) - Eq.(14) powers  $P_m$  from the location and characteristics of an orientation of a source is consequence of transition from Eq.(11) to Eq.(12). Physical preconditions of the given phenomenon are connected with repeated redeflection waves from walls block. This statement concerns only total sizes of powers, but does not concern to density of streams which, as shown further, pay off by means of the specified characteristics of a source.

Let's find ratio for calculation of good qualities  $Q_{vm}$ , by means of which factors  $g_m$  in Eq. (13) are defined.

Entered for the characteristic of absorbing properties of environment filling volume the parameter  $\beta$  is meaningful a tangent of a corner of losses dielectric with losses [17] owing to what partial good quality  $Q_{v0}$  is equal

$$Q_{\mu 0} = e^{-1}$$
 (15)

For calculation partial good qualities  $Q_{vm}$  at m> 1 we shall take advantage of a ratio

$$Q_{\mu m} = \operatorname{Re}(f_{\mu}) / \operatorname{Im}(2f_{\mu})$$
(16)

where  $f_{v}$ , - Eq.(4) frequency calculated on the formula at v-th a fashion provided that in examined structure there are no other losses of energy except for its outflow through m-that wall.

Walls surrounding structure are characterized by an impedance [17]

$$Z_{m} = c_{m} [1 - iy_{m} / (u_{\ell}e_{m})]^{-1/2},$$

where  $c_m = (e/M)^{1/2}$ ,  $\sigma_m$  and  $\varepsilon_m$  - characteristic resistance, electric conductivity and dielectric

 $\mathcal{U}_{\mathcal{H}}$ 

permeability of a material of which m-that wall is made,  $\mu$  - its magnetic susceptibility accepted identical to all materials in examined structure,  $\omega=2\pi f$ . Considering a ratio for depth of penetration of a field in a material of m-that wall [17]

$$\Delta_m = \frac{\sqrt{2}}{k_0} \left\{ \left| 1 + y_m^2 / (u_m e_m)^2 \right|^{1/2} - 1 \right\}^{-1/2}$$
(17)

where  $k_0 = u_i (e M)^{1/2} = 2p/\pi$ ,  $\pi$  - length of a wave in the environment filling volume of the block, the valid and imaginary parts of impedance  $Z_m$  can be written down in the form of:

$$\operatorname{Re} Z_{m} = \frac{c_{m} k_{0} \Delta_{m}}{\sqrt{2 + (k_{0} \Delta_{m})^{2}}} \cos u_{m}, \qquad (18)$$

Im 
$$Z_m = \frac{c_m k_0 \Delta_m}{\sqrt{2 + (k_0 \Delta_m)^2}} \sin u_{m,m}$$
  
where  $u_m = \frac{1}{2} \arctan\left(\frac{2}{k_0 \Delta_m} \sqrt{1 + (k_0 \Delta_m)^{-2}}\right)$ 

Named by a parameter of refraction of environment size  $n_m = c/c_m = (e_m/e)^{1/2}$  always it is more, and as a rule, it is much more 1, owing to what the maximal value of the module of impedance  $|Z_m|$  much less than characteristic resistance  $\rho$  the environment filling volume of examined structure  $\rho$ . For this reason value  $Z_m=0$ , m=1..., 6 it is used for calculation of initial approximation of resonant frequencies  $f_{\rm H}^0$  (see (4)) which, then, is specified by means of a method of indignation:

$$\operatorname{Re} f_{i} = f_{i}^{0} + \frac{\partial f_{i}}{\partial Z_{m}} \operatorname{Im} Z_{m} \approx f_{i}^{0}, \qquad (19)$$

$$\operatorname{Im} f_i = f_i^0 + \frac{\partial f_i}{\partial Z_m} \operatorname{Re} Z_m$$

where the derivative  $\partial f_{_{\rm H}}/\partial Z_m$  undertakes in point  $Z_1 = Z_2 = \ldots = Z_6 = 0$ .

In [18-19] the formula for a derivative  $\partial f_{\mu} / \partial Z_m$ 

$$\frac{\partial f_n}{\partial Z_m} = -\frac{1}{k_0 c k_{nm}^3} \frac{\partial f_n}{\partial h_m}$$
(20)

where  $h_m = V/S_m$  - the size of the block on a normal to m-th wall, V - its volume,  $S_m$  - the area of m-that wall is received,

$$k_{\mu m}^{\vartheta} = 1 - \int_{S_m} e \left| E_{\mu n} \right|^2 dS / \int_{S_m} \mathcal{M} \left| H_{\mu \phi} \right|^2 dS$$

 $E_{\text{Hn}}$ ,  $H_{\text{H}\varphi}$  - normal to m-that wall of a component electric and a tangent to it component of a magnetic field at v-th a fashion.

At  $Z_m=0$ , m=1..., 6 in examined structure degeneration E-and H-takes place aashions with an identical set of indexes n, l, and j. Believing as well as at transition from (11) to (12), that for frequencies from set  $\Omega$  of fashion E-and H-type with equal  $f_{\rm H}^0$  on the average bring the equal contribution to the energy saved up in the block and using formulas for components field Eand H-fashions, we receive [20-21]

$$\frac{1}{k_{_{HM}}^{_{9}}} = \frac{1}{\left(k_{_{HM}}^{_{9}}\right)^{E}} + \frac{1}{\left(k_{_{HM}}^{_{9}}\right)^{H}} = 1 + \frac{k_{_{0}}^{^{2}}}{k_{_{HM}}^{^{2}}} \quad (21)$$

where  $k_{\text{HM}}$ =op/h<sub>m</sub>,  $\xi$  - an index n, 1 or j, defining number of variations of a field of v-th fashion in a direction, perpendicular m-that wall.

Substituting Eq.(17) – Eq.(21) and following from Eq.(4) mode ratio

$$\frac{\partial f_i}{\partial h_{\dot{o}}} = -\frac{k_{i\dot{o}}^2}{k_0^2} \frac{f_i^0}{h_{\dot{o}}}$$

in Eq.(16) for m=1..., 6 it is found:

$$Q_{\mu m} = \left(1 + \frac{k_{\mu m}^2}{k_0^2}\right)^{-1} \frac{V}{2S_m \Delta'_m},$$
 (22)

where

$$\Delta'_m = \frac{n_m \Delta_m \cos \mu_m}{\sqrt{2 + (k_0 \Delta_m)^2}}$$

By means of ratio Eq.(13) – Eq.(15) and Eq.(22) it is possible to calculate components of losses of energy in the examined block. And, owing to presence of weight factor  $\partial_{\mu}$  in Eq.(12) satisfactory results on accuracy it is possible to expect even if in set of frequencies  $\Omega$  on which the required field decays, to include only fluctuations with  $|f_{\mu}^{0} - f_{\mu_{3,H}}|/f_{\mu_{3,H}} \leq 1$ .

After the termination of transients in each block into which the building is divided, there comes power balance for average powers: the total power which has arrived in any, for example, to k-th the block  $P_k^+$  it appears to the equal sum of powers disseminated in its volume  $P_k^0$  and leaving in surrounding walls  $P_k^{cm}$ :

$$P_{\kappa}^{+} = P_{\kappa}^{0} + P_{\kappa}^{cm} .$$
 (23)

And, power  $P_{k}^{cm}$  in Eq.(23) develops of powers  $P_{km}^{cm}$ , leaving in each of  $M_k$  surrounding the block walls

$$P_{\kappa}^{cm} = \sum_{m=1}^{M_{\kappa}} P_{\kappa m}^{cm} \,. \tag{24}$$

Power  $P_k^+$ , acting in k-th the block develops of the powers getting into it through walls from next blocks  $P_k^{cm}$  and, probably, powers of transmitter  $P_k^{H3\pi}$  if it is located in this block:

$$P_{\kappa}^{+} = P_{\kappa}^{cM} + 3_{\kappa} P_{\kappa}^{u3n}, \qquad (25)$$

where  $3_{\kappa}$  it is equal 1 for the block in which the transmitter is located, and 0 for all others.

Power  $P_{k}^{c_{M}}$  acting in k-th block from adjacent blocks with it, is directly proportional to size of powers following through their walls with the factor of easing depending on thickness and absorption in a material of which these walls are made.

Power  $P^{c_{M}}_{k}$  acting in to k-th the block from the next blocks through the common walls, it is possible to present in the form of:

$$P_{\kappa}^{cm} = \sum_{i=1}^{N} c_{ki} P_{i}^{cm}, \qquad (26)$$

where *N* it is equal to number of all blocks a building, elements  $c_{ki}$  are factors of easing of power  $P^{cm}_{i}$ , following through walls of i-th block and getting in to k-th the block,

$$c_{ki} = \begin{cases} 0, \\ \frac{S_{ki}}{S_{il}} g_{il} \exp\{-2d_{il} / \Delta_{il} \end{cases}$$

if blocks with numbers k and i do not have a common wall,

if blocks with numbers k and i have a common wall, (27)

 $S_{il}$  there is an area of l-st wall of i-th block,  $S_{ki}$  - the area overlapped with to k-th the block parts of this wall,  $g_{il}$  - factor with an index l for i-th block, counted on the formula Eq.(14).

Elements  $C_{ki}$  make a matrix C from dimension N × N which is meaningful matrixes of connections of blocks of a building. Its elements are defined by the plan of a building, and also thickness and absorbing properties of its external and internal walls and blockings. At absence of the common wall between to k-th and i-th blocks the factor  $C_{ki}$  is equal 0. Having substituted Eq.(23) – Eq.(24) and Eq.(26) in Eq.(25) it is received N the equations of a kind:

$$P_{\kappa}^{cm} = \left(1 - g_{k0}\right) \left[\sum_{i=1}^{N} c_{ki} P_{i}^{cm} + 3_{\kappa} P_{\kappa}^{u3n}\right], \qquad (28)$$
$$k = 1, \dots, N$$

Ratio Eq.(28) represent the equations of balance of powers in all blocks of a building and can be considered as system N of the linear algebraic equations rather N by unknown persons  $P_k^{cm}$ .

Having solved Eq.(28) for each block we find powers  $P_k^{cm}$ , by means of which then on Eq.(24) and Eq.(26) it is possible to find powers  $P_{km}^{cm}$  and  $P_{km}^{cm}$  the m leaving and acting in to k-th the block through its m-th wall. Having divided them on the area of corresponding walls  $S_m$ , we find resulting average density of streams of power  $T_{km}^{cp}$  at each wall in everyone block.

$$T_{km}^{cp} = \left(P_{km}^{0} + P_{km}^{cm}\right) / S_{m}.$$
 (29)

Stream  $T_{km}^{cp}$  is directed perpendicularly to m-that wall to k-th the block. Positive sign  $T_{km}^{cp}$  means, that power acts in to k-th the block through its m-that wall, negative-that through this wall power leaves from to k-th the block.

On known  $T_{km}^{cp}$  using procedure of interpolation it is possible to reconstruct size and direction of average density of a stream of power  $T^{cp}$  in any point inside any block of a building and further with its help to determine average accepted power in a vicinity of this point under the formula

$$P^{cp} = GT^{cp} \pi_{u_{3\pi}}^2 / (4p)$$
(30)

where G - factor of the directed action of the reception aerial,  $\Pi_{\mbox{\tiny H3JI}}$  - length of a wave of the transmitter.

In each block the law of change of accepted power at moving along any line has oscillatory character. In a vicinity of an any point of the block accepted power periodically accepts the maximal and minimal values caused by an interference of waves at their repeated redeflections. The formula Eq.(30) gives average value of this power in a vicinity of a an examined point of room. Having divided power of transmitter  $P^{II3I}$  on  $P_{np}^{cp}$ , we find average losses of distribution signal in each block of a building.

Repeated redeflection waves in each room lead to absence of any primary direction of distribution partial waves. Therefore, at resulting fluctuations in each room approximately it is possible to count identical all numbers  $k_{vm}$  in Eq.(21) – Eq.(22):

$$k_{\mu m}^2 = (1/3)(2p/\pi_{u 3 \pi})^2, \quad m = 1,...,M$$
 (31)

After substitution (31) in (22) and (14) it is received following simple formulas for calculation of factors  $g_{il}$  in (27):

$$g_{i0} = \frac{0.375\beta_i V_i}{0.375\beta_i V_i + \sum_{m=1}^6 S_{im} \Delta'_{im}},$$
 (32)

$$g_{il} = \frac{S_{il}\Delta'_{il}}{0.375\beta_i V_i + \sum_{m=1}^{6} S_{im}\Delta'_{im}} \quad (l \ge 1)$$

As it was already marked, through two - three walls from the location of the transmitter a determinative in accuracy of calculations is the size of accepted power  $P_{k0}^{cp}$  average on all examined (k-th) room. In this connection as a first approximation it is possible to do without procedure of interpolation and to calculate average power of an accepted signal in an any point of k-th room under the formula

$$p^{cp}(r) = P_{k0}^{cp}(r/r_{k0})^{-n}$$
(33)

where r - distance between the transmitter and the receiver,  $n = 1,5 \div 4$ ,  $r_{k0} = (r_k^{max} r_k^{min})^{1/2}$ ,  $r_k^{max}$  and  $r_k^{min}$ -distances from the transmitter up to the points most removed and closest to them to k-th a room (if the transmitter is located in a an examined room it is necessary  $r_k^{min} \approx 5\pi_{uan}$ ), and  $P_{\kappa o}^{cp}$  pays off on Eq.(30) for k-th room.

#### The conclusion

The developed model of communication channels inside of buildings will be realized in the form of the computer program and will find application for forecasting propagation losses of signals for frequencies of 900 and 1800 MHz in various types of uninhabited buildings: scientific and higher educational establishments, shopping centers. Among buildings to which approbation of model will be realized, can be multi-storey and one-storey, brick and concrete. Selecting the elementary block of separate rooms of a an examined building results of calculations of average signal propagation losses in each room should be compared to experimental data.

Thus it will be possible to estimate accuracy of the offered model.

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